



$$V_c(t) = \frac{1/pC}{1/pC + pL + R} V_{in}(t)$$

$$= \frac{1}{1 + p^2 LC + R p C} V_{in}(t)$$

$V_{in}(t) = A \cos \omega t \rightarrow$ use $\frac{A e^{j\omega t}}{2}$

$$V_c(t) = \text{Re} \frac{1}{1 - \omega^2 LC + R j \omega C} A e^{j\omega t}$$

take real

$$= \frac{1}{\sqrt{(1 - \omega^2 LC)^2 + (R \omega C)^2}} A \text{Re} (e^{j\omega t - j\phi})$$

$$\phi = \tan^{-1} \omega RC / (1 - \omega^2 LC)$$

$$\textcircled{1} V_c(t) = \frac{A}{((1 - \omega^2 LC)^2 + (R \omega C)^2)^{1/2}} \cos(\omega t - \phi)$$

Forced Response

Natural Response (Homogeneous solution)

This occurs when $V_c(t) \neq 0$ even if $V_{in}(t) = 0$

This can only happen when $1 + p^2 LC + R p C = 0$

or $p^2 + p \frac{R}{L} + \frac{1}{LC} = 0 \Rightarrow p = \frac{-\frac{R}{L} \pm \sqrt{\frac{1}{4}(\frac{R}{L})^2 - \frac{1}{LC}}}{2}$

= { $s_1 \rightarrow$ with + sign }
 { $s_2 \rightarrow$ with - sign }

$V_{ch} = B e^{s_1 t} + C e^{s_2 t}$ satisfies this.

since $\frac{1}{1 + p^2 LC + R p C} e^{s_1 t} = \frac{1}{1 + s_1^2 LC + R s_1 C} e^{s_1 t}$ and
 similarly for $e^{s_2 t}$

Total solution is the sum of

$$v_c(t) + v_{ch}(t) = v_{c\text{Tot}}(t)$$

Example: Find $v_c(t)$ when $v_c(0^+) = 0$
 $i_L(0^-) = 0$

when $v_{in}(t) = A \cos \omega t$

Solution

Since $v_c(0^+) = \frac{A}{((1 - \omega^2 LC)^2 + (R\omega C)^2)^{1/2}} \cos(\phi)$

$v_{ch}(0^+)$ must be $-v_c(0^+)$

Thus $\boxed{B + C = -v_c(0^+)}$ → need another condition (observe $i_c = C \frac{dv_c}{dt}$; $v_L = L \frac{di_c}{dt}$)

If v_L is finite $\frac{di_c}{dt}$ can change at $t=0$ BUT

$$i_c(0^+) - i_c(0^-) = \int_{0^-}^{0^+} \left(\frac{v_L}{L}\right) dt$$

→ 0 as 0^+ approaches 0^- . Thus

$i_c = i_L$ can't change instantaneously and $i_c(0^+) = 0$

$$i_c \Big|_{t=0} = C \frac{dv}{dt} = \frac{A}{((1 - \omega^2 LC)^2 + (R\omega C)^2)^{1/2}} (-\omega \sin \phi)$$

$$+ B s_1 + C s_2 = 0$$

Thus we determine B and C from

$$B + C = -\frac{A}{((1 - \omega^2 LC)^2 + (R\omega C)^2)^{1/2}} \cos \phi$$

$$B s_1 + C s_2 = +\frac{\omega A}{((1 - \omega^2 LC)^2 + (R\omega C)^2)^{1/2}} \sin \phi$$

We note that the homogeneous (or natural part of the solution) always decays because $p = -\frac{1}{2} \frac{R}{L}$ for s_1 and s_2

LT Spice plot of V_c

(3)

LT Spice plots the magnitude and phase of the steady-state! That is it plots (from page 1)

$$dB = 20 \lg \left[\frac{A}{\left((1 - \omega^2 LC)^2 + (R\omega C)^2 \right)^{1/2}} \right] / A$$

magnitude of $V_c(t)$ at low frequencies
 this eliminates A the excitation amplitude

$$= 20 \lg \left[\frac{\text{magnitude}(\omega)}{\text{magnitude}(0)} \right]$$

versus $\lg(\omega)$
 It also plots the phase ϕ

$$-\phi = -\tan^{-1} \omega RC / (1 - \omega^2 LC)$$

versus $\lg \omega$

we thus see $dB = 0$ for ω small

and $dB = 20 \lg \frac{1}{\omega^2 LC} = -40 \lg \omega - 20 \lg LC$ for ω large. It is a line of slope -40 dB/decade

Also ϕ starts at 0 for small ω and approaches $\phi \rightarrow -\tan^{-1} \infty = -\frac{\pi}{2}$ for $\omega = \frac{1}{\sqrt{LC}}$
 and $\phi \rightarrow -\tan^{-1} 0 = -\pi$ for $\omega \rightarrow \infty$

The height of the magnitude at $\omega = \frac{1}{\sqrt{LC}}$

(Eq 1) is given by $\frac{1}{R \frac{1}{\sqrt{LC}}} = \frac{1}{R \frac{\sqrt{LC}}{L}} = \frac{1}{R \omega_0 L} = \frac{\omega_0 L}{R}$

with $\omega_0 = \frac{1}{\sqrt{LC}}$. $\frac{\omega_0 L}{R} = Q$ is called the quality factor of the circuit