

# A tunable, ultrahigh spectral brightness KrF\* excimer laser source

R. T. Hawkins, H. Egger, J. Bokor,<sup>a)</sup> and C. K. Rhodes

Department of Physics, University of Illinois at Chicago Circle, P. O. Box 4348, Chicago, Illinois 60680

(Received 31 October 1979; accepted for publication 17 December 1979)

An extremely high spectral brightness KrF\* (248 nm) excimer source is described. This instrument combines the property of continuous tunability over the full gain profile with the following output pulse characteristics: pulse energy  $\sim 60$  mJ, pulse duration  $\sim 10$  nsec, spectral width  $150 \pm 30$  MHz, absolute frequency control to within 300 MHz, and beam divergence  $\sim 50$   $\mu$ rad. Within the uncertainty of measurement, the spectral width of the output radiation is Fourier transform limited, and the beam divergence corresponds to the diffraction of the radiating aperture.

PACS numbers: 42.60.By, 42.55.Hq, 07.65.Eh, 32.80.Kf

Rare-gas halide excimer sources have proven to be extremely useful in studies of multiphoton processes.<sup>1</sup> Nevertheless, the utility of these sources in such studies could be considerably enhanced. Heretofore, limitations of available sources have been (1) a broad ( $\sim 100$   $\text{cm}^{-1}$ ) emission profile, (2) the absence of a convenient, accurate, and reliable tuning system for control of the output wavelength, and (3) an output beam divergence on the order of one hundred times the diffraction limit. Overall, an enhancement of several orders of magnitude in spectral brightness, the key parameter in multi-quantum processes, is achievable if the output parameters of these sources are made to conform to the most stringent limits fundamentally possible.

Previous studies<sup>2,3</sup> have shown that substantial improvement in source properties can be obtained by using the output of a tunable, intracavity etalon line-narrowed, discharge-excited oscillator to injection lock a high-energy, unstable resonator oscillator. The best linewidth reported was  $\sim 10^{-1}$   $\text{cm}^{-1}$  in a beam with a divergence approximately twice the diffraction limit.<sup>3</sup>

In this letter, we report the properties of a KrF\* excimer source with performance parameters which closely approach the fundamental limits governing spectral width, beam divergence, and absolute wavelength control. The basic laser system is illustrated schematically in Fig. 1. The output of a frequency-stabilized, cw dye laser (Coherent 599-21,  $\Delta\nu < 5$  MHz,  $\sim 30$  mW at 497 nm) is pulse amplified in a 3-4 stage XeF\* pumped (30-50 mJ at 351 nm) dye amplifier, producing a 7-nsec visible pulse with an energy of  $\sim 3$  mJ at a repetition rate up to 1  $\text{sec}^{-1}$  (limited by thermal recovery time in the dye cuvettes). The linewidth of this visible radiation was examined with a scanning interferometer (Tropel 240) and was found to be  $85 \pm 10$  MHz. Frequency-doubled radiation corresponding to any wavelength within the KrF\* gain profile may readily be generated in a temperature-tuned, 90° phase-matched ADP crystal, producing  $\sim 5$  nsec second-harmonic pulses with energy  $> 100$   $\mu$ J. This spectrally narrow ( $120 \pm 20$  MHz FWHM, Fourier trans-

form limited) second-harmonic radiation is subsequently amplified in a single pass through a discharge-pumped KrF\* amplifier (Lambda Physik EMG 500) to produce output pulses of  $\sim 10$ -nsec duration, and energies up to 60 mJ.

A study of the spatial properties of the  $2 \times 0.5$ -cm-output beam reveal it to have a diffraction limited divergence of  $\sim 50$   $\mu$ rad. We note that at the peak measured power,  $\sim 6$  MW, focal intensities at 248 nm in excess of  $10^{15}$   $\text{W}/\text{cm}^2$  can be generated with f1 optics. In preliminary studies, we have observed optical breakdown in pure helium at pressures above 1.85 bars, with a KrF\* laser intensity of  $\sim 3 \times 10^{14}$   $\text{W}/\text{cm}^2$ .

Furthermore, since the frequency of the cw dye laser can be electronically scanned over 1  $\text{cm}^{-1}$ , continuously tunable coverage of the ultraviolet radiation over a 2- $\text{cm}^{-1}$  interval is readily accomplished. In addition, the dye laser frequency is known to within 300 MHz by interferometric comparison with a stable HeNe laser source.<sup>4</sup> In our current experiments, we have observed amplification at wavelengths  $\lambda_{\text{KrF}}$  in the range  $248.2 \text{ nm} < \lambda_{\text{KrF}} < 250.3 \text{ nm}$ . In this range, the output pulse energy was observed to monotonically decrease with increasing wavelength. In contrast to previous reports of a strong absorption centered at  $\sim 248.8$  nm in both *e*-beam pumped KrF\* amplifiers<sup>5</sup> and discharge-pumped tunable KrF\* lasers,<sup>2,6</sup> no reduction in output energy was observed at that wavelength in our apparatus.

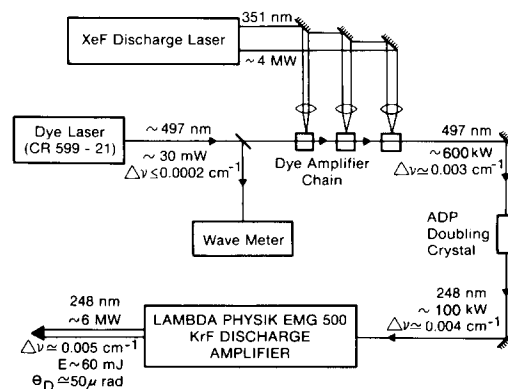


FIG. 1. Ultraviolet laser system illustrating the cw dye laser, the dye amplifier chain, the ADP doubling crystal, and the final KrF\* amplifier stage.

<sup>a)</sup>Also affiliated with the Department of Electrical Engineering, Stanford University, Stanford, California 94305.

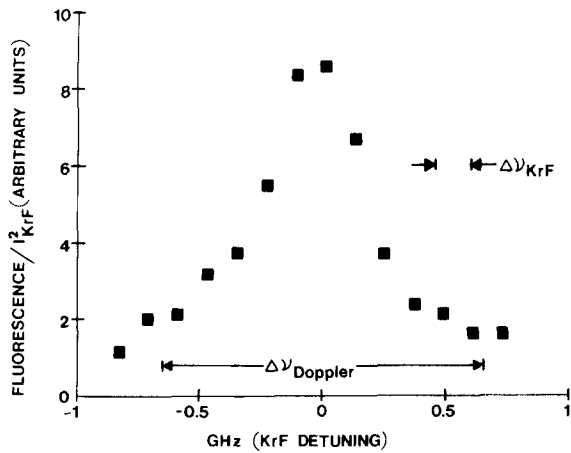


FIG. 2. Observed fluorescence at 828 nm following two-photon absorption in  $\text{Xe}(6p[\frac{1}{2}]_0 \leftarrow ^1S)$  at 249.63 nm, 0.25 Torr Xe pressure, and  $\sim 5 \times 10^5 \text{ W/cm}^2$  KrF\* intensity (unfocused).

In order to establish the linewidth of our source, direct interferometry has been used. A 1 GHz FSR Fabry-Perot interferometer was constructed with two dielectric coated mirrors of 98% reflectivity, resulting in a finesse of  $\sim 20$  for our typical 10-nsec pulse duration. With this interferometer, the amplified ultraviolet linewidth was measured to be  $\Delta\nu_{\text{KrF}} = 150 \pm 30 \text{ MHz}$ . We observe that this linewidth is equivalent, within our experimental uncertainty, to the Fourier-transform-limited linewidth  $\Delta\nu_{\text{SHG}} = 120 \pm 20 \text{ MHz}$  of the second harmonic output of the ADP crystal.

In order to further verify the narrow linewidth of the laser output, a preliminary study of the xenon ( $6p[\frac{1}{2}]_0 \leftarrow ^1S$ ) two-photon absorption at 249.63 nm was performed under Doppler-free conditions through detection of the resulting ( $6p[\frac{1}{2}]_0 \rightarrow 6s[\frac{3}{2}]_1$ ) fluorescence at 828 nm. Upon scanning the KrF\* source frequency across the resonance, a single feature was observed, centered at  $80\,118.73 \pm 0.10 \text{ cm}^{-1}$ , with a width of  $\sim 450 \text{ MHz}$  (FWHM, at the KrF\* frequency), as shown in Fig. 2. These measurements were made at a xenon pressure of 0.25 Torr and an ultraviolet intensity of  $\sim 5 \times 10^5 \text{ W/cm}^2$ . Although this width is a factor of 3 less

than the Doppler width, the resonance is still a factor of  $\sim 3$  broader than would be expected with a source linewidth of  $\sim 150 \text{ MHz}$ . Since we experimentally determined the pressure broadening coefficient of this xenon two-quantum transition to be  $25 \pm 15 \text{ MHz/Torr}$ , which gives a minor contribution to the linewidth at 0.25 Torr, we have concluded that unresolved isotope splittings (five xenon isotopes with natural abundances  $> 8\%$ ) lead to the observed width. An estimate<sup>7</sup> of the isotope shift for such a ( $5p^56p \leftarrow 5p^6$ ) transition yields a relative shift of  $\sim 50\text{--}100 \text{ MHz}$  for two Xe isotopes whose masses differ by 1 amu. This estimate agrees well with both the linewidth of the observed absorption, and the absence of additional resonances within  $\pm 6 \text{ GHz}$  of the one observed.

In conclusion, we state that an ultrahigh spectral brightness ultraviolet source, continuously tunable over the KrF\* excimer band, has been constructed with the following demonstrated performance characteristics:  $P \sim 6 \text{ MW}$ ,  $\Delta\nu_{\text{KrF}} = 150 \pm 30 \text{ MHz}$ , pulse repetition rate  $\sim 1 \text{ sec}^{-1}$ , and divergence  $\theta_D \simeq 50 \mu\text{rad}$ .

The authors wish to thank K. D. Skala and S. W. Vendetta for their technical assistance. This work was supported by the Office of Naval Research, the Department of Energy through agreement ED-78-S-08-1603, and the National Science Foundation through grant PHY78-27610. One of us (J.B.) also acknowledges support by the Fannie and John K. Hertz Foundation.

<sup>1</sup>W.K. Bischel, J. Bokor, D.J. Kligler, and C.K. Rhodes, *IEEE J. Quantum Electron.* **QE-15**, 380 (1979); C.K. Rhodes and P.W. Hoff in *Excimer Lasers*, edited by C.K. Rhodes (Springer-Verlag, Berlin, 1979), p. 175.

<sup>2</sup>J. Goldhar and J.R. Murray, *Opt. Lett.* **1**, 199 (1977).

<sup>3</sup>J.R. Murray, J. Goldhar, and A. Szöke, *Appl. Phys. Lett.* **32**, 551 (1978).

<sup>4</sup>F.V. Kowalski, R.T. Hawkins, and A.L. Schawlow, *J. Opt. Soc. Am.* **66**, 965 (1976).

<sup>5</sup>A.M. Hawryluk, J.A. Mangano, and J.H. Jacob, *Appl. Phys. Lett.* **31**, 164 (1977).

<sup>6</sup>T.R. Loree, K.B. Butterfield, and D.L. Barker, *Appl. Phys. Lett.* **32**, 171 (1978).

<sup>7</sup>I.I. Sobel'man, *Introduction to the Theory of Atomic Spectra*, (Pergamon, New York, 1972).