DOUBLE LAYER FORMATION IN A TWO REGION ELECTRONEGATIVE DISCHARGE

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OUTLINE

- Introduction
- "Experiments"
 - 2D particle-in-cell (PIC) simulations with rescaled oxygen cross sections
- Theory
 - 1D collisionless model of double layer (DL)
 - Global (volume-averaged) model of upstream and downstream particle and energy balances
- Comparisons
 - Reasonable agreement but also some differences
- Slow and fast wave instabilities
 - The DL has time-varying structure

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INTRODUCTION



• Why does a DL form at low pressures?

- The particle loss rate is greater upstream than downstream due to the smaller upstream radius
- A higher ionization rate (and T_e) is needed upstream than downstream
- A DL both "insulates" the low downstream T_e from the high upstream T_e , and it accelerates electrons upstream to increase the ionization rate there

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"EXPERIMENTS"

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PIC SIMULATION METHOD

- Self-consistent results from first principles with no assumptions about electron and ion velocity distributions
- Upstream heating at 13.56 MHz \perp to the plane of the simulation
- RF field adjusted to keep number of upstream electrons constant
- Stability, speed and accuracy require small plasma reactors with low densities and large Debye lengths:

 $n_e \approx 4 \times 10^{14} \text{ m}^{-3}, \lambda_D \approx 0.8 \text{ mm}$

- Due to low densities, rescaled oxygen cross sections were used: Positive-negative ion recombination $\uparrow \times 20$ Dissociative attachment $\uparrow \times 5$
- A typical simulation takes 1–2 weeks
- The pressure range explored is 0.5-24 mTorr
- $\bullet\,$ No DL was observed at 0.5 mTorr

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PIC RESULTS FOR 6 mTorr O₂ DISCHARGE



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RESULTS FOR 6 mTorr O₂ DISCHARGE (CONT'D)



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THEORY

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DL MODEL

Potential $V = V_s$ **D**

Upstream

edge 2

- Charge density ρ and potential V within the DL are found by solving Poisson's equation
- Six types of particles contribute to ρ:

 thermal positive ions, negative ions, and electrons
 accelerated positive ions, negative ions, and electrons
- The particle motions are 1D and collisionless
- The boundary conditions are that ρ and $d\rho/dV$ vanish at the DL edges
- An additional condition is that the sum of positive and negative charge in the double layer must vanish; equivalently, the total force acting on the double layer must vanish

$$\int_0^{V_s} \rho(V) \, dV = 0$$

• A final condition that upstream and downstream-directed electron fluxes nearly balance determines the equilibrium value of V_s/T_h .

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Downstream

edge 1 V=0



- We use a 2D rectangular geometry
- To determine the equilibrium quantities we use
 - global particle balance upstream
 - global particle balance downstream
 - global energy balance downstream
- Upstream energy balance (which determines the upstream electron density) depends on the input power, and is beyond the scope of this study

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COMPARISONS

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MODEL (TOP) and PIC (BOTTOM) EEDF'S (6 mTorr O_2)







TIME-VARYING PHENOMENA

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SLOW AND FAST WAVE INSTABILITIES (6 mTorr O_2)



- At 2–12 mTorr, the DL coexists with an unstable slow wave that originates downstream and propagates upstream as it grows
- The wave frequency is 50–100 kHz with a wavelength of order 1 cm
- The wave produces $\sim 20\%$ oscillations in the double layer potential and ~ 0.5 cm oscillations in the DL position
- We believe the wave is driven by counter-streaming flows of positive and negative ions (Tuszewski and Gary, 2003)
- There is an additional unstable fast wave at higher pressures having frequency of order 1 MHz and a wavelength of order 1 cm LiebermanISPC09

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MOVIE SHOWING SLOW AND FAST WAVES

Red solid line: fast waves averaged over 0.1475 μ s intervals Blue solid line: slow waves averaged over 1.18 μ s intervals

160. Waves in 6 mTorr DL region (23.6 microsecs)



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CONCLUSIONS

- 2D PIC simulations can be powerful tools to study the physics of double layer formation
- The simulations provide diagnostics at the edges and inside the DL, which would be very difficult to do in a laboratory experiment
- The simulations show upstream edge Maxwellian electrons along with an accelerated component, and a downstream edge bi-Maxwellian with a hotter tail tied to the upstream temperature
- The simulations show that the DL can coexist with unstable slow and fast wave-driven oscillations
- A 1D analytical DL model coupled with global models of the upstream and downstream particle and energy balance captures the essential physics of DL formation in this two-region system

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